Confirmation of the 1^{-+} Meson Exotics in the $\eta\pi^0$ System

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(Dated: June 21, 2007)

The exclusive reaction $\pi^-p \to \eta\pi^0 n$, $\eta \to \pi^+\pi^-\pi^0$ at 18 GeV/c has been studied with a partial wave analysis on a sample of 23 492 $\eta\pi^0 n$ events from BNL experiment E852. A mass-dependent fit is consistent with a resonant hypothesis for the P_+ wave, thus providing evidence for a neutral exotic meson with $J^{PC} = 1^{-+}$, a mass of $1257 \pm 20 \pm 25$ MeV/c^2, and a width of $354 \pm 64 \pm 60$ MeV/c^2. New interpretations of the meson exotics in neutral $\eta\pi^0$ system observed in E852 and Crystal Barrel experiments are discussed.

PACS numbers: 13.20.Jx; 13.85-t; 14.40.Cs

Keywords: meson spectroscopy; exotic mesons

INTRODUCTION.

Exotic mesons with $J^{PC} = 0^{--}, 1^{-+}, 2^{+-}, \ldots$ do not mix with quark-antiquark mesons and thus offer a natural testing ground for QCD. Exotic mesons have been discussed [1–10] for many years but have only recently been observed experimentally. The underlying structure of the negatively charged exotic state with $J^{PC} = 1^{-+}$ observed in this experiment [11, 12] at 1400 MeV decaying into $\eta\pi^-$ is not yet understood.

Two distinguishing characteristics of the $\eta\pi^0$ system make it an excellent one to clarify the properties of this exotic state. First, C-parity is a good quantum number in the $\eta\pi^0$ system, unlike the $\eta\pi^-$ system. Second, the production mechanism for the charge exchange reaction $\pi^-p \to \eta\pi^0 n$ cannot involve the exchange of an isospin $I = 0$ system such as the pomeron.

The Crystal Barrel experiment [13] confirmed the existence of resonant structure in the $\eta\pi^-$ system using stopped antiprotons in liquid deuterium in the reaction $\bar{p}n \to \pi^-\pi^0\eta$. Later this group analyzed data on $\bar{p}p$ annihilation at rest into $\pi^0\pi^0\eta$ [14] and presented evidence for an exotic 1^{-+} resonance in the $\eta\pi^0$ system with $M = (1360 \pm 25)$ MeV/c^2 and $\Gamma = (220 \pm 90)$ MeV/c^2.

The $\eta\pi^0$ state has been studied in the GAMS experiment [15] with $\pi^-p \to \eta\pi^0 n$, $\eta \to 2\gamma$, $\pi^0 \to 2\gamma$ at 32, 38 and 100 GeV/c. The statistics of the 38 GeV/c data was sufficient so that, using the method of Sadovsky [16], they were able to present evidence for the exotic $\pi_1(1400)$.

The VES experiment also observed a peak in the P_+ wave of the $\eta\pi^0$ system near 1400 MeV/c^2 [17]. In their most recent publication [18], using theoretical arguments the authors state that the peak can be understood without requiring an exotic meson.

An analysis of E852 data with $\pi^-p \to \eta\pi^0 p$, $\eta \to 2\gamma$ was recently reported [19]. A bump in the P_+ wave of the $\eta\pi^0$ system was observed at $M(\eta\pi^0) = 1272$ MeV/c^2 with a large width ($\Gamma = 660$ MeV/c^2) when fitting all the data using the method of Ref. [12]. For small t', they observe a width of 190 MeV/c^2. The authors chose not to claim evidence for exotic $\pi_1(1400)$ meson production.

The present analysis studies $\pi^- p \to \eta\pi^0 n$, $\eta \to \pi^+\pi^-\pi^0$ at 18 GeV/c in E852. The advantage of this mode over the all-neutral final state is that the production vertex is defined by charged tracks. This improves the mass resolution as well as the ability to require that the interaction took place within the target.

EXPERIMENTAL SETUP AND DATA SELECTION

The data for this analysis were obtained at the Alternating Gradient Synchrotron (BNL USA) in 1995. Using an 18 GeV/c π^- beam interacting in a liquid hydrogen target, a total of 750 million triggers were acquired of which 108 million were of a type designed to enrich the exclusive final state $\pi^- p \to \pi^+\pi^- 4\gamma n$. A total of 6 million events of this type were fully reconstructed. The data were kinematically fit [20] to select events consistent with the $\pi^-\pi^+\pi^0\pi^0 n$ hypothesis (with a confidence level of at least 0.01%) yielding about 4 million events. Of those, 85 228 events passed a mass cut enhancing η mesons, $m(\pi^-\pi^+\pi^0) < 0.65$ GeV/c^2, and 74 549 passed a cut to remove events passing through a low-efficiency region in the drift chambers. A final kinematic fit selected 23 492 events for the partial wave analysis (PWA), which were consistent with the $\eta\pi^0 n$, $\eta \to \pi^+\pi^-\pi^0$ hypothesis at a minimum confidence level of 1%.

A strong η meson signal is observed in this final data sample (Fig. 1a) with a mass of 539.2 ± 0.3 MeV/c^2 and a (resolution-dominated) width of 23.7 ± 0.2 MeV/c^2. The filled regions in the figure indicate the signal region and the side-band regions used in the analysis. In the η signal region, the signal-to-background ratio is about 4 to 1 for all $\eta\pi^0$ masses and 5 to 1 for $m(\eta\pi^0) > 1.1$ GeV/c^2. The $\eta\pi^0$ mass spectrum shown in Fig.1b has two clear peaks: the $a_0^0(980)$ and the $a_2^0(1320)$.

The non-η background was estimated as a function of $\eta\pi^0$ mass using the side-band and signal regions. The background fraction varies between 24% and 14% going from lower to higher mass in the region $0.78 < m(\eta\pi^0) < 1.74$ GeV/c^2.

The experimental acceptance was determined using a Monte Carlo event sample generated with isotropic angular distributions in the Gottfried-Jackson frame. The detector simulation was based on the E852 detector simulation package SAGEN [12]. The experimental acceptance was incorporated into the PWA by means of Monte Carlo normalization integrals [12]. The acceptance as a function of mass and as a function of t' is flat.

PARTIAL WAVE ANALYSIS

The partial-wave analysis (PWA) method described in [12] (see also [21, 22]) was used to study the spin-parity structure of the $\eta\pi^0$ system in this data set. The PWA was carried out using the extended maximum likelihood method separately in each mass bin in the mass region between 0.78 and 1.74 GeV in mass bins of 0.04 GeV for $0 < |t'| < 1.0$ (GeV/c)^2 using the likelihood function

$$\ln\mathcal{L} \propto \sum_i^n \ln I(\Omega_i) - \int d\Omega \eta(\Omega)\, I(\Omega). \tag{1}$$

Here $I(\Omega)$ is the predicted angular distribution, $\eta(\Omega)$ is the angular acceptance, and the sum is over the event sample.

The partial waves are parameterized by a set of five numbers: $J^{PC}m^\epsilon$, where J is the angular momentum, P and C are the parity and the C-parity of the $\eta\pi^0$ system, m is the absolute value of the angular momentum projection and ϵ is the reflectivity. We use a simplified notation where each partial wave is denoted by a letter indicating the $\eta\pi^0$ system's angular momentum in standard spectroscopic notation, and a subscript which can take the values 0, +, or −, for $m^\epsilon = 0^-$, 1^+, or 1^- respectively. We assume that the contribution from partial waves with $m > 1$ is small and can be neglected [12].

The amplitudes used are the unnatural parity-exchange waves (UNPW) S_0, P_0, P_-, D_0, D_-, and the natural parity-exchange waves (NPW) P_+, D_+. The NPW waves interfere between themselves as do the UNPW waves but the NPW waves do not interfere with the UNPW waves. The P_+ wave would be an exotic $J^{PC} = 1^{-+}$ (denoted by π_1^0) if the wave is resonant.

For each partial wave the complex production amplitudes were determined from an extended maximum likelihood fit [22]. The spin 1/2 nature of the target proton leads to spin-flip and spin-non-flip amplitudes and thus to a production spin-density matrix with maximal rank two. The PWA fit presented in this paper was carried out with the assumption that a spin-density matrix of rank one was sufficient [12]. An isotropic incoherent background was included. The magnitude of the background was fixed as determined from the side bands. We investigate the quality of the fits by comparing the moments of the decay angular distributions $H(LM), L \leq 4$ [12, 22], of the data with those predicted

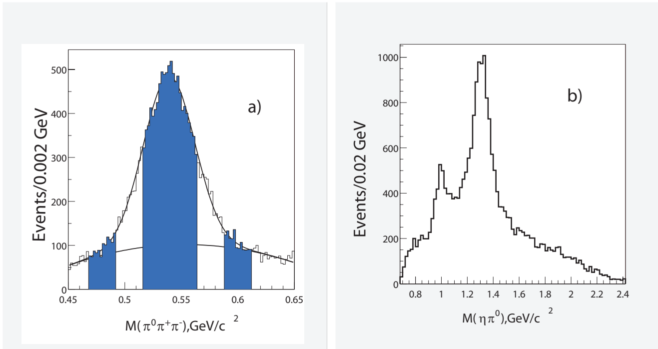

FIG. 1: (a) Fit of the $\pi^+\pi^-\pi^0$ mass distribution in the η mass region. There are two entries per event: one for each way to assign a π^0 to the η decay. The η signal region and the side-band regions are shown shaded. (b) The uncorrected $\eta\pi^0$ effective-mass distribution for events consistent with the reaction $\pi^- p \to \eta\pi^0 n$

by Monte Carlo events generated with the fit amplitudes. We also directly compare the angular distributions for $\cos(\theta_{GJ})$ and φ_{TJ} between the data and those Monte Carlo events. The quality of the fits is good.

Since natural-parity exchange (NPE) and unnatural-parity exchange (UNPE) amplitudes have different $|t'|$ dependences, a fit to the $|t'|$ distribution using a function of the form $N(t') = n_1|t'|e^{-b_1|t'|} + n_2 e^{-b_2|t'|}$ was carried out to determine the relative contributions of the two exchanges. The fitted parameters are $b_1 = (7.41 \pm 0.08)(\mathrm{GeV/c})^2$, $b_2 = (2.68 \pm 0.07)(\mathrm{GeV/c})^2$, and thus the ratio of UNPE and NPE contributions is equal to 0.71 ± 0.03. A value of about 70% for the ratio of UNPE to NPE at 18 GeV/c is expected based on interpolation of experimental $\pi^- p$ charge exchange data between 3 and 40 GeV/c in the Regge model [16].

There is mathematical ambiguity in the description of a system of two pseudo-scalar mesons [23]. For our set of amplitudes there are eight ambiguous solutions, each of which leads to identical angular distributions. These solutions were found analytically starting from one solution by means of the Barrelet zeros method [22]. The eight solutions in each mass bin are shown in Fig. 2 as a point for every ambiguous solution. (In some cases the solutions are too close together to be visible as separate solutions. This is particularly true in the mass bins at $M = 1.56$ GeV/c^2 and at $M = 1.72$ GeV/c^2.) The spread between the various ambiguous solutions dominates the systematic uncertainty in the resonance parameter determination (see below).

MASS DEPENDENT FITS

To study the resonant structure in the partial waves, we used three different procedures. The first method utilizes a Mass Dependent Fit (MDF) of the average solutions in the NPW sector. The PWA results in each mass bin were averaged between ambiguous solutions [12]. The mass dependence of the P_+ and D_+ intensities as well as their relative phase difference were then fit by relativistic Breit-Wigner (BW) functions (in both the P_+ and D_+ waves) with mass-dependent widths and Blatt-Weisskopf barrier factors[12]. The mass and width of the a_2^0 are well known and were fixed using values[24]: $M = 1320$ MeV/c^2, and $\Gamma = 120$ MeV/c^2. The width of the $a_2(1320)$ includes the experimental mass resolution. There are three free parameters in the fit of the D_+ intensity, $|D_+|^2$: one for the magnitude and two parameterizing the smooth background for the D_+ wave, as was done in [12].

In the MDF of the $|P_+|^2$ distribution and the relative phase $\Delta\Phi(D_+ - P_+)$ there are four free parameters: three from the BW function and one for the production phase (assumed constant). The fit was carried out in the mass interval $1.1 - 1.74$ GeV/c^2. The resonant hypothesis for D_+ and P_+ waves with a mass-independent production phase gives a χ^2/DoF=1.14 for 28 degrees of freedom. The non-resonant hypothesis (no phase variation for the P_+ wave) gives χ^2/DoF=3.02. It is clear from Fig. 2c that a single resonant phase for the $a_2(1320)$ (dotted line) with a constant (non-resonant) P_+ wave is not satisfactory.

The contrast between the $D_+ - P_+$ phase variation and the D_+ phase variation with mass (Fig. 2c, solid line) shows clearly that the observed phase variation is consistent with interference between two resonant waves. It's worth pointing out that the mass dependence of the $D_+ - P_+$ relative phase would be nearly flat if each wave were resonant with very similar masses and widths.

The P_+ resonant parameters from the fit with the average solutions and the average error matrix [12] are: $M =$

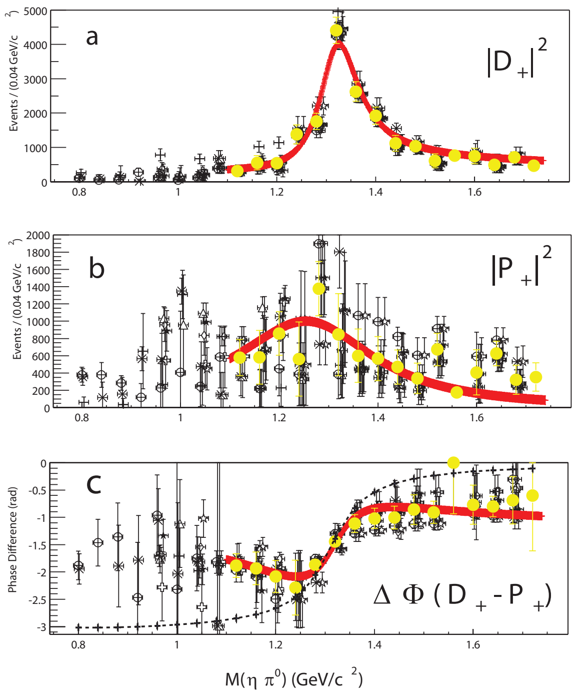

FIG. 2: The Partial Wave Analysis (PWA) and Mass-Dependent Fit (MDF) results. The points shown in each mass bin are the eight ambiguous PWA solutions. a) the D_+ wave intensity; b) the P_+ wave intensity; and c) the relative phase between the P_+ and D_+ waves. The lines show the MDF results (method 2). The average PWA solution in each mass bin is plotted using grey (yellow) points. The dotted line in (c) is the phase difference if the P_+ phase is constant.

1265 ± 21 MeV/c^2 and $\Gamma = 411 \pm 64$ MeV/c^2.

The second method was similar to the first except that instead of fitting the average solutions, a large number ($\simeq 10^3$) of randomly chosen combinations of ambiguous solutions in each mass bin were used as input to the mass-dependent fit. The obtained distributions of the mass and width of the P_+ resonance for those fits with acceptable values of χ^2 (that is, those with $\chi^2/DoF < 2$) were then fitted by a Gaussian. The mean values of these distributions are: $M = 1257 \pm 20$ MeV/c^2 and $\Gamma = 354 \pm 64$ MeV/c^2. The curves shown in Fig. 2 are drawn using these mean values.

The RMS values of these distributions are $\sigma_M = 25$ MeV/c^2 and $\sigma_\Gamma = 58$ MeV/c^2. The systematic errors are obtained using these RMS values. Our analysis shows that this spread due to the different ambiguous solutions dominates the systematic error.

The third method used was to carry out a Mass-Dependent Partial Wave Analysis (MDPWA) [12]. In this procedure, an extended maximum likelihood function is generalized to include not only the angular distribution, but also the $\eta\pi^0$ mass distribution for each wave. This analysis is free from the problem of ambiguous solutions but it is necessary to parameterize the mass dependence of every partial wave (including the UNPWs) and all relative phases. We use the same parametrization for the D_+ and P_+ waves as in the first two methods. The mass dependence of the UNPW waves were chosen to be polynomials of second order with constant phases except for the S_0 wave. The S_0 wave was fitted with a BW function using the $a_0(980)$ resonance parameters. The MDPWA results for the P_+ wave are $M = 1256 \pm 10$ MeV/c^2 and $\Gamma = 319 \pm 34$ MeV/c^2, consistent with the results of methods 1 and 2.

In [12] it was shown that a pure D_+ wave can artificially induce a P_+ wave due to the effects of finite acceptance and resolution. This "leakage" leads to a P_+ wave that mimics the D_+ intensity and phase. In our case, the P_+ intensity would therefore have an intensity with the shape of the $a_2(1320)$ and a $D_+ - P_+$ phase difference which doesn't depend on mass. These features allowed us to include in the MDPWA fit a leakage term with these features. We observed that the leakage contribution to the P_+ wave from the D_+ wave is negligible.

Evidence for a resonance interpretation for the P_+ wave is primarily the behavior of the $D_+ - P_+$ relative phase (Fig. 2c). Since the D_+ phase variation is well known because of the $a_2(1320)$ production, it is clear that the P_+ phase cannot be constant (see dotted curve in Fig. 2c) and it is well-described by a BW phase variation.

The ratio of the P_+ and D_+ intensities in the range $1.24 < M(\eta\pi^0) < 1.34$ GeV is equal to $|P_+|^2/|D_+|^2 = 0.43 \pm 0.10$.

This ratio is larger than that for the $\eta\pi^-$ system, as reported in Ref. [12].

The mass of the neutral exotic 1^{-+} state, decaying into $\eta\pi^0$, observed here ($M = 1257 \pm 20 \pm 25$ MeV/c^2) is lower than the mass observed in the Crystal Barrel experiment [14] ($M = 1360 \pm 25$ MeV/c^2) by about 100 MeV although the results are barely consistent within errors. The width measured here ($\Gamma = 354 \pm 64 \pm 58$ MeV/c^2) is also consistent with that from the Crystal Barrel measurement ($\Gamma = 220 \pm 90$ MeV/c^2).

It should also be noted that our result is similar to those obtained from the low-t' fits in Ref. [19]. They obtained $M = 1301 \pm 14$ MeV/c^2 and $\Gamma = 190 \pm 32$ MeV/c^2 in one analysis, and $M = 1386 \pm 32$ MeV/c^2 and $\Gamma = 363 \pm 81$ MeV/c^2 in a fit to the experimental moments. No systematic errors were given for either method.

The lower mass found in our analysis may be a consequence of interference between the resonant state and background in the $\eta\pi^0$ system, some of which may be from rescattering between the η and the π^0. It is also possible that two or more resonant 1^{-+} states may be present in the $\eta\pi^0$ decay channel in the mass interval mass between 1200 and 1400 MeV as might be expected if the exotic state is a four quark state.

The mass for the π_1^0 is some 100 MeV below the negatively charged π_1^- measured by Crystal Barrel [13] and BNL-E852 [11, 12]. Of course, this does not constitute compelling evidence for more than one exotic state in this mass region. However, if we *assume* that there are two $\eta\pi$ states at 1280 and at 1380 MeV respectively, then one expects the lower-mass state to have a substantial branching ratio into $\rho\pi$ (the reaction $\pi^- + p \to \eta\pi^0 + n$ is mediated by ρ exchange), whereas the higher-mass state would couple to the $f_2(1270)\pi$ channel or perhaps to the $Pomeron + \pi$ channel (the reaction $\pi^- + p \to \eta\pi^- + p$ is produced in part by the exchange of the Pomeron).

CONCLUSIONS

Mass dependent fits of the D_+ and P_+ amplitudes and their relative phase using three different methods (described above) all lead to the conclusion that the P_+ wave is well-described by a resonance hypothesis and is inconsistent with having a constant phase. The resonance parameters for the observed π_1^0 are given by $M = 1257 \pm 20 \pm 25$ MeV/c^2 and $\Gamma = 354 \pm 64 \pm 58$ MeV/c^2. Here the first error is statistical and the second is systematic. (We have chosen to take the resonant parameters and errors from method 2.) This result, together with the previous results from Crystal Barrel [14] and E852 [19] provide strong indications for one or more spin-exotic mesons near 1400 MeV/c^2 decaying to $\eta\pi^0$.

We would like to express our appreciation to the members of the MPS group, and to the staffs of the AGS, BNL, and the collaborating institutions for their efforts. This research was supported in part by the National Science Foundation, the US Department of Energy, and the Russian Ministry for Education and Science. The Southeastern Universities Research Association (SURA) operates the Thomas Jefferson National Accelerator Facility for the United States Department of Energy under contract DE-AC05-84ER40150.

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